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THE MAGNETIC FIELD SIGN REVERSAL
AND EVOLUTION OF ROSSBY AND ALFVÉN WAVES
INDUCED BY VELOCITY SHEAR OSCILLATIONS ON THE SUN

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The behaviour of the toroidal and meridional components of the solar large-scale magnetic field and linear Alfvén and Rossby waves during solar activity cycles and bi-annual time periods are theoretically investigated in this work. The shear of the toroidal and meridional components of velocity is taken into account. We consider the case of periodical velocity shear with bi-annual oscillation period, hereinafter the Velocity Shear Quasi Bi-annual Oscillations (VSQBO). The large-scale magnetic field toroidal and meridional components are obtained as harmonic functions of the time and the oscillation phase difference between them is equal to $\pi/2$. The sign reversal of the magnetic field toroidal and meridional components is studied. The numerical simulations show that in the case of significant VSQBO amplitude values the toroidal or meridional component of the large-scale magnetic field reverses its sign three times in one of the hemispheres (northern or southern) of the Sun, during the solar activity cycle 23. According to our results the appearance of velocity shear oscillations leads to the modulation of the magnetic field 22-year period oscillations by the bi-annual ones. The presented model is applicable for investigation of the magnetic field evolution at the base of convection zone as well as for understanding the magnetic field properties in the upper solar atmosphere (the magnetic canopy of the solar chromosphere Roberts (1996)).

The excitement of the linear Rossby and Alfvén waves in the shear layer at the base of the convection zone is also considered. The periodical impulsive growth is characteristic to the energy density of Alfvén and Rossby waves and they propagate as localized in time powerful pulses. Properties of these pulses are governed by the direction of waves phase velocity and the initial phase of VSQBO, and such behavior of waves well explains mechanisms of the solar flare excitement and activity.
1. Introduction

The large-scale magnetic field longitudinal and latitudinal distribution and sign reversal during the solar cycle 23 have been investigated by many authors (de Toma et al., 2000; Benevolenskaya et al., 1999; Benevolenskaya, 1996; Benevolenskaya and Makarov, 1992; Bai, 1995; Bumba, 1989; Gaizauskas et al., 1983; Vitinskij, 1982; Ambroz, 1973). The recent observations of the p-mode frequency splitting and results of corresponding helioseismic inversions of the solar internal rotation show the appearance of two different layers with strongly marked shear flow (e.g. Kosovichev and Schou, 1997; Sekii, 1998). One of these layers is located at the base of convection zone and the differential rotation in this region, as is well known, plays an important role in the solar dynamo mechanism. The second one is located just below the photosphere. It is also reasonable to think that layer effects on the distribution of the global magnetic field at the solar surface and takes place in the formation of canopy of magnetic field in the upper solar atmosphere Roberts (1996). In other words the differential rotation takes an imposing part in the formation and evolution of the large-scale solar magnetic field.

In this work we present the ideal magnetohydrodynamic (MHD) model explaining three-time sign reversal of the large-scale magnetic field at the beginning of cycle 23. It is known from the theory of the solar dynamo Parker (1979), that the toroidal and meridional components of the large-scale magnetic field oscillate with constant phase difference of $\pi/2$: the toroidal magnetic field reaches its maximum value when the meridional one is minimum or vice versa. In our model we obtain toroidal and meridional magnetic field components as harmonic functions of the time, when the differential rotation and its time oscillation are taken into account. Hereinafter in the work, the periodic time variation of the velocity shear with 2-year period is denoted as the Velocity Shear Quasi Bi-annual Oscillations (VSQBO). We present the result of numerical simulations showing that our model can explain one-time and three-time sign reversals of the large-scale magnetic field components during the activity cycle 23 and, in general, the modulation of activity cycle 23 by oscillations with a period of about 2 years. The behaviour of the large-scale magnetic field strongly depends on the amplitude and initial phase of VSQBO. The similar model was developed by Pataraya et al. (1997) using VSQBO. In that case the toroidal component of the magnetic field was a linear function of time and the meridional one was a constant. The structure of the large-scale magnetic field cyclic oscillations is also considered in the several works (Benevolenskaya, 1994, 1995, 1998; Akioka et al., 1987).

We analytically investigate the possible mechanisms of the formation and evolution of the active regions and the corresponding excitement of the solar flares. In this sense, excitement of solar flares and their periodical activity are well known from several observations (e.g. see Bai, 1987; Rieger et al., 1984). In this work we investigate the propagation of the linear Rossby and Alfvén waves in the environment with the steady and periodically oscillating velocity shear. Additionally, toroidal and meridional components of the magnetic field change sign one or three times during the activity cycle 23. We consider the time evolution of these waves. The numerical simulations show that the hydrodynamic and the magnetic energy densities oscillate with short periods and the waves propagate as localized in time pulses. These results are in agreement with the observed periodicity of solar flare activity and this allows one to think that our model is applicable for the explanation of solar flares excitement and activity.

2. Basic Equations

In this work we investigate the behaviour of the large-scale solar magnetic field in the layer, with the differential rotation and excitement of linear waves in this layer. As it is evident from helioseismic observations there are two kinds of such layers in the solar interior. One of them is at the base of the convection zone and the other is just below the photosphere. These layers represent the areas bounded by two concentric spherical surfaces. The calculations are done in the co-rotating Cartesian frame. $xoy$ plane is tangential to the bottom boundary of the layer.
The $x$ axes is directed eastward and it represents the toroidal direction, the $y$ axes is directed northward and it represents the meridional direction and the $z$ axes is directed outward to the upper solar atmosphere and it determines the radial direction. We use the ideal MHD equations for incompressible medium (Priest, 1982). So, for given latitude we have a set of differential equations in the following form:

$$\frac{DB_x}{Dt} = B_x \frac{\partial V_x}{\partial x} + B_y \frac{\partial V_x}{\partial y}, \quad (1)$$

$$\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} = 0, \quad (2)$$

$$\frac{DV_x}{Dt} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + \frac{B_x}{4 \pi \rho} \left( \frac{\partial B_y}{\partial y} - \frac{\partial B_y}{\partial x} \right) + (f_0 + \beta y)V_y, \quad (3)$$

$$\frac{DV_y}{Dt} = -\frac{1}{\rho} \frac{\partial p}{\partial y} + \frac{B_x}{4 \pi \rho} \left( \frac{\partial B_y}{\partial y} - \frac{\partial B_y}{\partial x} \right) - (f_0 + \beta y)V_x, \quad (4)$$

$$\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} = 0, \quad (5)$$

where,

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + V_x \frac{\partial}{\partial x} + V_y \frac{\partial}{\partial y}. \quad (6)$$

In these equations $V_x, B_x$ are the toroidal and $V_y, B_y$ are the meridional components of the velocity and the magnetic field, respectively; $\rho = \text{const}$ and $p$ are density and pressure; $f_0 = 2\Omega \sin \theta$ is Coriolis parameter;

$$\beta = \frac{2}{R_\odot} \frac{\partial}{\partial \theta} (\Omega \sin \theta), \quad (7)$$

where, $\Omega$ is the angular rotation frequency of the sun; $R_\odot$ is the Solar radius and $\theta$ is the latitude corresponding to the origin of the local coordinate frame. The positive values of $\theta$ correspond to the northern hemisphere and the negative ones denote latitudes in the southern hemisphere. The other symbols have their usual meanings.

It is well known from the solar models that the rotation frequency of the Sun is a function of latitude $\Omega = \Omega(\theta)$ and it could be well approximated as (Brown et al., 1989; Howard and Harvey, 1970):

$$\Omega = \Omega_0 (1 - A_1 \sin^2 \theta - A_2 \sin^4 \theta), \quad (8)$$

where, $\Omega_0, A_1$ and $A_2$ are constants. It should be noted that by using eq. (7) we obtain magnitude of the latitude, where $\beta$ is equal to zero. After substituting $A_1$ and $A_2$ values we have $\beta = 0$ at the following latitudes:

$$\theta_1^{(N)} \approx 74^{\circ}33', \quad \theta_1^{(S)} \approx -74^{\circ}33' \quad (9)$$

As we have mentioned above our aim is to study the excitement of the linear Rossby and Alfvén waves in the layer where steady velocity shear and VSQBO appear. But before we turn to this matter let us determine the mathematical form of quantities satisfying the (1)-(5) set of equations in an unperturbed state of environment.

### 3. The Sign Reversal of the Magnetic Field Components

We consider the differential rotation of the Sun with respect to the co-rotating Cartesian frame of reference described above. In this case we have the toroidal component of the unperturbed velocity in the following form:

$$V_{x0} = gV_{xy}, \quad (10)$$
where, $V_{xy}$ is the shear in the $y$ direction of the toroidal flow and it has the form:

$$V_{xy} = V + V_{xn} \cos(\omega_n t + \varphi_0) = V_f^2,$$  \hspace{1cm} (11)

The first term on the right side of eq. (11), $V$ is the permanent shear of the toroidal flow in the considered layer and it conveys the steady differential rotation of the Sun:

$$V = \frac{d}{d\theta}(\Omega \cos \theta) = -\Omega_0 \sin \theta(1 + 2A_1 + (4A_2 - 3A_1) \sin^2 \theta - 5A_2 \sin^4 \theta).$$  \hspace{1cm} (12)

This expression was presented in Pataraya and Zaqlarashi (1995) and Gigolashvili et al. (1995). [In these works the excitement of linear waves were investigated, when the toroidal component of the unperturbed magnetic field is a linear function of time.]

By using eq. (12) we can obtain the latitude value, where the steady velocity shear of the toroidal flow $V$ has a maximum by solving equation $dV/d\theta = 0$, and after substituting of $A_1$ and $A_2$ values we obtain:

$$\theta^{(N)}_2 \approx 55^\circ 9', \quad \theta^{(S)}_2 \approx -55^\circ 9'$$  \hspace{1cm} (13)

The other term on the right side of eq. (11) determines the velocity shear oscillation with an arbitrary period and in the case of 2 year oscillation period it demonstrates the VSQBO. In this term $\omega_n$ is the frequency of the velocity shear oscillations:

$$\omega_n = \frac{2\pi}{nt_0},$$  \hspace{1cm} (14)

where $t_0 \approx 3.16 \cdot 10^7 \text{sec} = 1 \text{ year}$. $nt_0$ is the oscillation period of the velocity shear.

We have the meridional component of the unperturbed velocity in the form:

$$V_{y0} = xV_{yx},$$  \hspace{1cm} (15)

where $V_{yx}$ is the shear of the meridional flow in the $x$ direction:

$$V_{yx} = -\frac{\Omega_0^2}{V} f_2,$$  \hspace{1cm} (16)

According to eqs. (10)-(12), (15)-(16) and the equation of induction we obtain expressions for the toroidal $B_{x0}$ and meridional $B_{y0}$ components of the unperturbed magnetic field:

$$B_{x0} = B_{xm} \sin(\Omega_1 t_0 f_1),$$  \hspace{1cm} (17)

$$B_{y0} = B_{ym} \cos(\Omega_1 t_0 f_1),$$  \hspace{1cm} (18)

where, $\Omega_1 = 2\pi/22t_0$ is the frequency of low-frequency oscillations (with the period 22 years) of the unperturbed toroidal and meridional components of the solar magnetic field; $B_{xm} = B_0 \sin \theta = V \cdot B_{ym}/\Omega_1$ here, $B_0 = \text{const}$

$$f_1 = \tau - 1 + a_n[\sin\left(\frac{2\pi}{n} \tau + \varphi_0\right) - \sin(\varphi_0)],$$  \hspace{1cm} (19)

$$f_2 = 1 + \frac{V_{xn}}{V} \cos\left(\frac{2\pi}{n} \tau + \varphi_0\right),$$  \hspace{1cm} (20)

where $\tau = t/t_0$ and

$$a_n = \frac{V_{xm}}{\omega_n t_0 V}.$$  \hspace{1cm} (21)

As we can see from eqs. (17) and (18) the toroidal and meridional components of the magnetic field are harmonic functions of the time and they oscillate with a constant phase difference of $\pi/2$. We carried out the numerical simulations for the case of $n = 2$ in eq. (14). In other words, we
consider the VSQBO. We studied the time evolution of the magnetic field components defining the VSQBO amplitude in two different ways:

a) Using in eqs. (11) and (21) the following expression of the \( V_{xn} \):

\[
V_{xn} = V_0 \sin^2 \theta \cos \theta,
\]

one has

\[
a_2 = -\frac{V_0}{\pi \Omega_0} \frac{\sin \theta \cos \theta}{1 + 2A_1 + (4A_2 - 3A_1) \sin^2 \theta - 5A_2 \sin^4 \theta}.
\]

Analysing the results of numerical simulations we found that the low-frequency (with 22-year period) oscillation of the large-scale magnetic field components is modulated by the high-frequency (with bi-annual period) one. The modulation rate depends on the latitude \( \theta \) and the magnitude of \( a_2 \). Particularly, the significant modulation of the magnetic field oscillation takes place, when \( \pi |a_2| > 1 \) and the three-time sign reversal of the magnetic field toroidal component occurs at the beginning of the activity cycle. Figure 1 shows the time evolution of the normalized toroidal \( B_x/B_{xm} \) (solid lines) and the meridional \( B_y/B_{ym} \) (dashed lines) components. Figure 1a corresponds to the latitude value, where the steady velocity shear of the toroidal flow has maximum \( \theta = \theta_2^{(N)} \) in the northern hemisphere and Figure 1b corresponds to the same latitude in the southern hemisphere (see eq. (13)). In this case we assume that the amplitude of the VSQBO \( V_{xn} \) is defined by eq. (22) and its initial phase \( \varphi_0 = \pi \). These results show that at the middle latitudes the amplitude of VSQBO is large enough to induce the significant modulation of the large-scale magnetic field cyclic oscillations by the bi-annual ones in both hemispheres of the Sun. This modulation leads to the magnetic field toroidal component three-time sign reversal within the first two years of the activity cycle 23 in the northern hemisphere (see the encircled area 1 in Figure 1a), while in the southern one only one-time sign reversal occurs during the same time interval. As one can see from the presented Figures, when \( \varphi_0 = \pi \) at the solar activity minimum the three-time sign reversal of the magnetic field toroidal component appears only in the northern hemisphere. We have a similar picture in the southern hemisphere in the time interval between 11 and 13 years from the beginning of the cycle 23. The time behaviour of the magnetic field toroidal component is inverted within this time interval, i.e. now three-time sign reversal occurs only in the southern hemisphere (see the encircled area 1 in Figure 1b). In other words, the modulation of the magnetic field components oscillation is not synchronized between different hemispheres and, because of this, we have the asymmetry in the three-time sign reversal occurrence. This kind of behaviour is characteristic of the magnetic field toroidal component for most latitudes except the areas very close to the equatorial plane and pole. We will address this in more detail later. With regard to the magnetic field meridional component (the dashed line curves), its low-frequency oscillation is also strongly modulated by the bi-annual one in both hemispheres of the Sun. In spite of this, the three-time sign reversal of the meridional field component disappears for most of the latitude values, although this effect can almost be achieved only close to the latitudes where the steady velocity shear is maximum (eq. (13)) and this effect occurs in the northern hemisphere in the time interval \( 6 < \tau < 8 \) years (see the encircled area 2 in Figure 1a) and in the southern one in the time interval \( 5 < \tau < 7 \) years (see the encircled area 2 in Figure 1b). It is important to note that, if we perform calculations for the case of \( \varphi_0 = 0 \), then we obtain an inversed picture, i.e., in this case the three-time sign reversal of the magnetic field toroidal component occurs only in the southern hemisphere, within the first two years of activity cycle.

As we mentioned above the situation changes in the equatorial and polar areas. The time evolution of the magnetic field components in these areas is very similar in both hemispheres, we therefore present the numerical results regarding the northern hemisphere. In Figures 1c and 1d the curves for latitudes close to the equatorial plane (\( \theta = 5^\circ \)) and to the pole (\( \theta = 85^\circ \)) in the northern hemisphere are presented, respectively. The behaviour of the large-scale magnetic field in the equatorial and polar areas is different from that considered above. In these areas the amplitude of the VSQBO is small and, because of this, \( \pi |a_2| \) becomes < 1. As one can see in
Figures 1c and 1d, in this case the very weak modulation of the magnetic field low-frequency oscillation occurs and the above mentioned sign reversal does not appear. Finally, exactly on the equatorial plane \((\theta = 0^\circ)\) or on the poles \((\theta = \pm 90^\circ)\) the modulation of the cyclic oscillation vanishes completely.

b) We have also performed the numerical simulations for the case when the amplitude of the VSQBO \(V_{xn}\) is defined as:

\[
V_{xn} = -U_0 V \sin(\theta(1 - (\omega_n t_0)^{-1}) + \sin \theta), \tag{24}
\]

and in this case we obtain \(a_2\) as follows:

\[
a_2 = -\frac{U_0 \sin \theta(\delta(1 - \pi^{-1}) + \sin \theta)}{\pi}, \tag{25}
\]

where \(U_0\) and \(\delta\) are constants. As in the first case, we present in Figure 2 the curves showing the time dependence of the normalized toroidal and meridional magnetic field components, when the VSQBO amplitude is defined by eq. (24), its initial phase is \(\varphi_0 = \pi\), \(\delta = 1.4\) and \(U_0 = 1.8\). Figure 2a corresponds to the latitude value in the northern hemisphere, \(\theta = \theta_2^{(N)}\), and Figure 2b to that in the southern one, \(\theta = \theta_2^{(S)}\). These results show, that unlike the previous case, when \(\delta = 1.4\) the strongly marked modulation of the magnetic field component low-frequency oscillation appears (i.e. \(\pi|a_2| > 1\)) only in the northern hemisphere, while in the southern one \(\pi|a_2| < 1\) and the modulation is weak. Hence, the three-time sign reversal of the magnetic field toroidal component within first two years of the activity cycle occurs only in the northern hemisphere (see the encircled area 1 in Figure 2a), while in the southern hemisphere the modulation is very weak and thus the three-time sign reversal is absent in the whole hemisphere. Unlike the previous case, the asymmetry in the magnetic field toroidal component three-time sign reversal is not due to different phases of modulation of the low-frequency oscillations, but now the strong modulation occurs only in the northern hemisphere. As it is shown in Figure 2a the oscillation of the meridional component is modulated too, but in spite of this its three-time sign reversal does not occur at most of latitude values. It is important to note that, we can obtain an inversed picture in the case, when \(\delta = -1.4\) and \(\varphi_0 = 0\). i.e. in this case we have the strong modulation of the magnetic field cyclic oscillation only in the southern hemisphere.

Finally, the behaviour of the magnetic field components close to the equatorial plane is similar to that shown in the previous case. As one can see from Figure 2c, the above mentioned modulation is very weak and thus the sign reversal effect does not occur near the equatorial plane. Although, from eq. (25) one can see that \(|a_2|\) increases with latitude \(\theta\) and when \(\delta > 0\) it is maximum on the northern pole and when \(\delta < 0\) it is maximum on the southern pole. Hence, in these cases the three-time sign reversal the magnetic field toroidal component appears at one of the polar areas too (see the encircled area 1 in Figure 2d). Moreover, in this case the three-time sign reversal of the meridional component can also be achieved (see the encircled area 2 in Figure 2d).

4. The Evolution of Rossby and Alfvén Waves

In the previous section we studied (1)-(5) a set of differential equations in the unperturbed state and we determined the expressions for the toroidal and meridional components of the unperturbed flow velocity and magnetic field. Let us now consider the excitement of the linear waves in the incompressible medium located in the solar internal shear layer. In the unperturbed state we take into account the shear and VSQBO (eqs. (10) and (15)). Additionally, the components of the magnetic field are oscillating with period of 22 years. At the same time this low-frequency oscillation is modulated by the high-frequency oscillation of the magnetic field with period equal to that of VSQBO.
We investigated equations (1)-(5) by using the linear perturbation theory. So, all quantities are presented in the form $\psi = \psi_0 + \psi'$ and we obtain the following set of the linearized differential equations:

$$\frac{\partial B'_x}{\partial x} + \frac{\partial B'_y}{\partial y} = 0,$$

(26)

$$\frac{\partial V'_x}{\partial x} + \frac{\partial V'_y}{\partial y} = 0,$$

(27)

$$\frac{D_0 B'_x}{D_0 t} = \left( B_{x0} \frac{\partial}{\partial x} + B_{y0} \frac{\partial}{\partial y} \right) V'_y + V_{yx} B'_x,$$

(28)

$$\frac{D_0}{D_0 t} \left( \frac{\partial V'_x}{\partial y} - \frac{\partial V'_y}{\partial x} \right) = \frac{1}{4 \pi \rho} \left( B_{x0} \frac{\partial}{\partial x} + B_{y0} \frac{\partial}{\partial y} \right) \left( \frac{\partial B'_x}{\partial y} - \frac{\partial B'_y}{\partial x} \right) + \beta V'_y,$$

(29)

where,

$$\frac{D_0}{D_0 t} = \frac{\partial}{\partial t} + V_{x0} \frac{\partial}{\partial x} + V_{y0} \frac{\partial}{\partial y}.$$

(30)

As we can see from eqs. (10) - (15) both $V_{x0}$ and $V_{y0}$ depend on the time and spatial coordinates. On the other hand the magnetic field components are harmonic functions of time (eqs. 17-18). Because of this, in the equations we have temporal as well as the spatial inhomogeneity. In this case the modal analysis is not applicable for the investigation of the perturbed quantities.

We studied eqs. (26)-(29) by using the non-modal perturbation theory which is considered by many authors (Chagelishvili et al., 1993; Lominadze et al., 1988; Farrell, 1987; Marcus and Press, 1977; Taylor, 1931). According to non-modal analysis, when we consider the linear velocity shear profile, it is possible to find out the coordinate transformation, which in the linearized equations, leads to the changing of spatial inhomogeneity by a temporal one. Using this coordinate transformation allows one to investigate the temporal evolution of the spatial Fourier harmonics and in this case all of the perturbed functions have the following form:

$$\psi' = \psi_1 \exp\{iK_xx + iK_yy\},$$

(31)

where,

$$K_x = k_{x1} F_1,$$

(32)

$$K_y = \frac{V}{\Omega_1} k_{x1} F_2,$$

(33)

are the toroidal and meridional components of the wavevector, respectively. In expressions (32) and (33):

$$F_1 = C K_{21} \sin(\Omega_1 t_0 f_1) + \cos(\Omega_1 t_0 f_1)$$

(34)

$$F_2 = - \sin(\Omega_0 f_1) + C K_{21} \cos(\Omega_0 f_1),$$

(35)

where, $C = \Omega_1/V$ and $K_{21} = k_{y1}/k_{x1}$ In these equations $k_{x1}$ and $k_{y1}$ are constants. According to expressions (26)-(29) and (31)-(33) we obtain the set of equations for perturbed functions in the following form:

$$K_x B_{x1} + K_y B_{y1} = 0,$$

(36)

$$K_x V_{x1} + K_y V_{y1} = 0,$$

(37)

$$\frac{dB_{y1}}{dt} = i \sqrt{4 \pi \rho \omega A} V_{y1} - \frac{K_y}{K_x} V_{yx} B_{y1},$$

(38)

$$\frac{d}{dt} \left( \frac{K_x^2}{K_y} V_{y1} \right) = i \beta V_{y1} + i \omega A \frac{K_x^2}{\sqrt{4 \pi \rho} K_x} B_{y1},$$

(39)

where,

$$K_x^2 = K_{x1}^2 + K_{y1}^2.$$

(40)
and $\omega_A$ is the frequency of Alfvén waves and it has the following form:

$$\omega_A = \frac{K_x B_{x0} + K_y B_{y0}}{\sqrt{4\pi\rho}} = \frac{B_{zm}}{\sqrt{4\pi\rho}} C k_y 1,$n)

(41)

where, $B_{zm}$ is the amplitude of the magnetic field toroidal component (see eq. (17)). First, we consider the case when $k_y 1 = 0$. In this case $\omega_A = 0$ and propagation of only Rossby waves is possible and we obtain the analytical solution $V_{y1} = V_{y1}^{(R)}$ of eq. (39) in the following form:

$$V_{y1}^{(R)} = V_{y1}(0) \frac{\cos(\Omega_1 t_0 f_1)}{1 + (C^{-2} - 1) \sin^2(\Omega_1 t_0 f_1)} \exp\left\{ \frac{i\beta}{k_x 1 (C^{-2} - 1)^{1/2}} \arctg(\sqrt{(-2 - 1)^{1/2} \sin(\Omega_1 t_0 f_1)}) \right\}$$

(42)

Since all perturbed quantities are assumed to be $\sim \exp(i\varphi)$ and the oscillation frequency $\omega = -\partial \varphi/\partial t$ we obtain the dispersion relation of Rossby waves:

$$\omega_R = -\frac{i\beta \cos(\Omega_1 t_0 f_1)}{k_x 1 (1 + (C^{-2} - 1) \sin^2(\Omega_1 t_0 f_1))} + x k_x 1 \Omega_1 f_2 \sin(\Omega_1 t_0 f_1) + y k_x 1 \Omega_1 f_2 \cos(\Omega_1 t_0 f_1),$$

(43)

and at last, by using of eqs. (37) and (42) we have the expression of Rossby waves hydrodynamic energy density $W_R = \rho(|V_{x1}^{(R)}|^2 + |V_{y1}^{(R)}|^2)/2$ in the following form:

$$W_R = \frac{\rho |V_{y1}^{(R)}|^2(0)}{2} \frac{1}{1 + (C^{-2} - 1) \sin^2(\Omega_1 t_0 f_1)}$$

(44)

In Figure 3 the time dependence of the $W^R$ normalized by its initial value is presented. Figures 3a and 3b correspond to the northern ($\theta = \theta_2^{(N)}$) and southern ($\theta = \theta_2^{(S)}$) hemispheres, respectively, when the amplitude of the VSQBO is defined by eq. (22) and its initial phase $\varphi_0 = \pi$. As we can see from these figures, when in the medium propagate only Rossby waves their hydrodynamic energy periodically increases in short time period during the solar activity cycle, which is demonstrated by peaks in the plots. The time moments of the wave energy density peak appearance coincide with those when the unperturbed magnetic field toroidal component changes its sign. As we have mentioned above, in this case occurs three-time sign reversal of the toroidal component and this fact reveals itself in the three peaks at the beginning of the activity cycle 23 (Figure 3a). The peaks denoted by "1" in Figures 3a and 3b represent the the moment of sign reversal due to the mean low-frequency oscillations of the magnetic field toroidal component and the occurrence of peaks denoted by "2" and "3" is caused by the three-time sign reversal effect. Hence, when one-time sign reversal occurs one has only the single peaks. In the moments of sign reversal of the unperturbed magnetic field toroidal component the Rossby waves propagate as pulses localized in time. [It should be noted that by comparison of Figure 3a and 3b it is clearly seen that the asymmetry between the northern and southern hemispheres due to the fact that modulation of the magnetic field cyclic oscillation in the different hemispheres is asynchronous.]

We have also performed the numerical simulations in the case when the amplitude $V_{zn}$ of the VSQBO is defined by eq. (24), its initial phase $\varphi_0 = \pi$, $\delta = 1.4$ and $U_0 = 1.8$. Resulted curves presented in Figure 4 showing the time dependence of the Rossby wave hydrodynamic energy density normalized by its initial value. Figures 4a and 4b correspond to the northern ($\theta = \theta_2^{(S)}$) and southern ($\theta = \theta_2^{(S)}$) hemispheres, respectively. In these Figures we have the peaks of energy density in the time moments when sign reversal of the unperturbed magnetic field toroidal component occurs. As we know from the above discussion in this case we have three-time sign reversal only in the northern hemisphere and this leads to the appearance of two additional peaks denoted by "2" and "3" beside the mean peak denoted by "1" in Figure 4a. Furthermore, we have only single peaks in the southern hemisphere and this is a result of the disappearance of the strong modulation of the magnetic field cyclic oscillations by bi-annual
ones. This property of waves determines the asymmetry of solar flare activity intensity within
the solar hemispheres and a detailed discussion on this matter is presented below.

It is important to note that, in general, in the case of \( k_y = 0 \), the meridional component
of the Rossby wave wavevector is \( K_y = 0 \) at the time moments when the sign reversal of the
unperturbed magnetic field toroidal component occurs. This means that, in time moments when
the Rossby waves propagate exactly in the toroidal direction the wave energy density increases
and it has the form of a sharp peak.

Now let us consider the case when \( k_y \neq 0 \). In this case, in the medium, propagation of both
Rossby and Alfvén waves is possible. We performed a numerical calculation of the (36)-(39) set of
differential equations. In Figure 5 the normalized hydrodynamic and magnetic energy densities
are presented as functions of time corresponding to \( \theta = \theta^{(N)}_2 \) (Figures 5a) and \( \theta = \theta^{(S)}_2 \) (Figures
5b) values of the latitude. One can see from these curves that wave hydrodynamic

\[
W_H = \frac{\rho}{2} (|V_{x1}|^2 + |V_{y1}|^2)
\]

and magnetic

\[
W_M = \frac{|B_{z1}|^2 + |B_{y1}|^2}{8\pi}
\]

energy densities periodically increase during the short time intervals while in the initial mo-
mement we assumed that the hydrodynamic and magnetic energy densities are commensurable
\( W_H(0)/W_M(0) \approx 1 \). The linear Alfvén waves propagate as powerful pulses. The period of energy
density oscillation depends on the \( k_{x1} \) and \( k_{y1} \) magnitude. In other words the hydrodynamic
and magnetic wave energy densities oscillate with period which depends on the direction of
the wave phase velocity in the \( \tau = 0 \) time moment. Figure 5 shows curves for \( k_{x1} = 10^{-9} \)
cm\(^{-1} \) and \( K_{z1} = 2 \). The calculations are carried out for values of environment parameters
satisfying those obtained from the standard solar models for the base of the solar convection zone.
Particularly, the amplitude value of the unperturbed magnetic field toroidal component (eq. (17))
is determined assuming \( B_0 = 10^5 \) G. Figure 6 shows the time variation of the hydrodynamic and
the magnetic energy densities for the same values of parameters, but in this case the amplitude
of the VSQBO is taken by eq. (24). As it is clear from the curves in Figure 5 and Figure 6 the
character of the wave temporal evolution is similar in both cases. This model well explains the
mechanism of the solar flares excitement and activity. For the different values of \( k_{x1} \) and \( k_{y1} \) the
wave energy density oscillation with periods well describing the observed 51-day (e.g. see Bai,
1987) as well as 154-day (e.g. see Rieger et al., 1984) periodicity in the occurrence of the solar
flare activity may be obtained. Observational data (Bai, 1987) also show asymmetry in intensity
of the flare activity between northern and southern hemispheres of the Sun. Comparing curves
demonstrated in Figures 5a and 5b (Figures 6a and 6b) one can see that the amplitude of energy
density oscillation in northern hemisphere is greater than that in southern hemisphere and the
more intensive excitement of the solar flares in the northern hemisphere is expected. According
to our results the flare activity is more intensive in the hemisphere where the three-time sign
reversal of the unperturbed magnetic field toroidal component occurs within the first two years
of the activity cycle.

As we have shown in section 2 in the unperturbed state the modulation of large-scale magnetic
field low-frequency oscillation vanishes close to the polar and equatorial areas, when the VSQBO
amplitude is determined by eq. (22). This leads to the decrease of the wave energy oscillation
amplitude in mentioned areas. On the other hand, when we consider the case of the VSQBO
amplitude by eq. (24) then in the polar area of the northern hemisphere the amplitude of the
wave energy density oscillations remains to be significant.

So, we have presented the model of solar flare excitement at the base of the solar convection
zone as the reasonable mechanism of the solar flare activity origin. As we mentioned in the
introduction our model is also applicable for the shear layer just below the photosphere. It is
reasonable to think that the velocity shear effects on the large-scale magnetic field distribution
at the solar surface and the excitement and the time evolution of Rossby and Alfvén waves take
part in the formation of the magnetic field structure in the upper solar atmosphere and in the
heating of the solar corona.

5. Conclusions

In this work we investigated the behaviour of the large-scale magnetic field in the solar internal
shear layers. The VSQBO is taken into account. By use of the numerical simulations we studied
the behaviour of the magnetic field toroidal and meridional components. The excitement and
time evolution of the linear Rossby and Alfvén waves at the base of convection zone were also
studied, when the VSQBO and modulation of the magnetic field components were taken into
consideration. We obtained results for our analytical solutions and for the above mentioned
numerical simulations, which can be formulated as follows:

1) We obtained the expressions of the magnetic field toroidal and meridional components in
the form of harmonical functions of time.

2) We obtained the curves representing the time dependence of the magnetic field toroidal
and meridional components in two cases:

a) When the VSQBO amplitude is defined by eq. (22). In this case, as shown in Figures 1a
and 1b, we have the strongly marked modulation of the large-scale magnetic field low-frequency
(with period 22-years) oscillation by the high-frequency (with the 2-year period) ones at most
latitudes of both northern and southern hemispheres of the Sun. When the initial phase of
VSQBO \( \varphi_0 = \pi \), the modulation of the magnetic field components cyclic oscillation leads
to three-time sign reversal of the toroidal component in the northern hemisphere, within the first
two years after activity minimum. During the same time interval we have only one-time sign
reversal of the magnetic field toroidal component in the southern hemisphere. This asymmetry,
in occurrence of the three-time sign reversal effect between different hemispheres, is caused
by asynchronous modulation phase in these hemispheres. The three-time sign reversal of the
magnetic field meridional component does not happen at most latitude values. Exceptions are
areas where the steady shear of the velocity toroidal component is maximal (see eq. (13)). In
this case the three-time sign reversal of the meridional component may happen in time interval
between 6 and 8 years after activity minimum in the northern hemisphere and between 5 and 7
years in the southern hemisphere.

The above mentioned effect of the modulation decreases in equatorial and polar areas (Figures
1c and 1d). Because of this, in these areas three-time sign reversal of any magnetic field
component becomes impossible.

b) We also considered another case, when the VSQBO amplitude is defined by eq. (24).
Corresponding curves are presented in Figure 2. In this case, when the initial phase of VSQBO
\( \varphi_0 = \pi \) the strongly marked modulation of the magnetic field components low-frequency oscil-
lation appears only in the northern hemisphere (see Figure 2a), while in the southern hemisphere
this modulation is very weak (Figure 2b). Because of this the three-time sign reversal of the
field toroidal component occurs only in the northern hemisphere, within the first two years after
activity minimum. This asymmetry is the result of the modulation range difference between
solar hemispheres.

Similar to the previous case the modulation range decreases in equatorial area, where the
three-time sign reversal effect vanishes. But in this case the magnetic field cyclic oscillation
modulation range is maximal in the northern polar area. Hence, in this area modulation remains
significant. Moreover, near the northern pole the three-time sign reversal of the magnetic field
meridional component becomes possible.

Further, we considered the excitement of the linear Rossby and Alfvén waves in the base
of convection zone, when in unperturbed state VSQBO and corresponding modulation of the
magnetic field cyclic oscillations are taken into account. We carried out the analytical and
numerical investigation of the excitement and evolution in time of these waves and we obtained the following results:

3) When in eqs. (34)-(35) \( k_{y1} = 0 \), i.e. \( K_{21} = 0 \) in the environment propagate only Rossby waves. In this case it is possible to analytically solve the set of differential equations (36)-(39) and we have obtained a solution of the meridional component of the perturbed velocity (see eq. (42)). By using this solution we determined the expression of the Rossby wave hydrodynamical energy density temporal variation.

4) We obtained the curves of the Rossby wave energy density time dependence. Curves in Figure 3 correspond to the definition of the VSQBO amplitude by eq. (22) and Figure 4 corresponds to that defined by eq. (24). These figures show that the Rossby wave energy density increases during very short time intervals and such behaviour of the wave energy density are shown as sharp peaks. These peaks appear in the time moments when the sign reversal of the unperturbed magnetic field toroidal component occurs. Thus, in time moments when the three-time sign reversal of the magnetic field toroidal component happens we have three peaks instead of one. In the case of \( k_{y1} = 0 \) these time moments coincide with those of the wavevector meridional component \( K_y = 0 \). This means that in these moments Rossby waves propagate exactly in the toroidal direction.

5) We performed the numerical investigation of the set of differential equations (36)-(39) in case when \( k_{y1} \neq 0 \). In this case propagation of both Rossby and Alfvén waves is possible. We obtained the curves of the wave hydrodynamic and magnetic energy density temporal variation. Figure 5 corresponds to the case "a" of the VSQBO amplitude and Figure 6 corresponds to case "b". These results clearly show oscillation of the wave hydrodynamic and magnetic energy densities. By choosing the \( k_{x1} \) and \( k_{y1} \) values one can obtain oscillation of wave hydrodynamical and magnetic energy densities with periods comparable with those obtained from observations of the solar flare activity (e.g. see Bai, 1987; Rieger et al., 1984). These reasons allow us to think that the model we have presented here well describes the solar flare excitement and their periodical activity. Our results are in agreement with the observations of solar flare activity which show the difference of the flare activity intensity between northern and southern hemispheres of the Sun (Bai, 1987).

The model presented here is applicable for investigation of the magnetic field variations and the linear wave evolution in the layers where the arbitrary period oscillations of the velocity shear, characteristic to Solar atmosphere, occur. In this sense it is important to study the role of the shear layer just below the photosphere in the distribution of the magnetic field at the solar surface. This model is also useful for the investigation of the excitement and evolution of waves close to the solar surface and their influence on the solar corona. Finally, taking into account the velocity shear oscillation with 11-year period, these work results could be used for investigation of solar p-mode properties during activity cycles. These very important tasks will be subjects of our forthcoming works.

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Figure Captions

Figure 1. The normalized toroidal $B_{x0}/B_{xm}$ (solid line) and meridional $B_{y0}/B_{ym}$ (dashed line) components of the solar large-scale magnetic field are plotted against $\tau$, time (Years). The VSQBO amplitude is defined by eq. (22) and its initial phase is $\varphi_0 = \pi$. a) the latitude $\theta = \theta_2^{(N)}$; b) $\theta = \theta_2^{(S)}$; c) the equatorial area $\theta = 5^\circ$; d) the polar area $\theta = 85^\circ$.

Figure 2. The normalized toroidal $B_{x0}/B_{xm}$ (solid line) and meridional $B_{y0}/B_{ym}$ (dashed line) components of the solar large-scale magnetic field are plotted against $\tau$, time (Years). The VSQBO amplitude is defined by eq. (24) and its initial phase is $\varphi_0 = \pi$. a) the latitude $\theta = \theta_2^{(N)}$; b) $\theta = \theta_2^{(S)}$; c) the equatorial area $\theta = 5^\circ$; d) the polar area $\theta = 85^\circ$.

Figure 3. The curve of the normalized Rossby wave hydrodynamical energy density $W_R(\tau)/W_R(0)$ as a function of $\tau$, time (Years). The VSQBO amplitude is defined by eq. (22) and its initial phase is $\varphi_0 = \pi$. a) the latitude $\theta = \theta_2^{(N)}$; b) $\theta = \theta_2^{(S)}$.

Figure 4. The curve of the normalized Rossby wave hydrodynamical energy density $W_R(\tau)/W_R(0)$ as a function of $\tau$, time (Years). The VSQBO amplitude is defined by eq. (24) and its initial phase is $\varphi_0 = \pi$. a) the latitude $\theta = \theta_2^{(N)}$; b) $\theta = \theta_2^{(S)}$.

Figure 5. The normalized Alfvén wave hydrodynamical $W_H(\tau)/W_H(0)$ (solid line) and the magnetic $W_M(\tau)/W_M(0)$ (dashed line) energy densities are plotted against $\tau$, time (Years), when in the initial momentum $W_H(0)/W_M(0) \approx 1$. The VSQBO amplitude is defined by eq. (22) and its initial phase is $\varphi_0 = \pi$. a) the latitude $\theta = \theta_2^{(N)}$; b) $\theta = \theta_2^{(S)}$.

Figure 6. The normalized Alfvén wave hydrodynamical $W_H(\tau)/W_H(0)$ (solid line) and the magnetic $W_M(\tau)/W_M(0)$ (dashed line) energy densities are plotted against $\tau$, time (Years), when in the initial momentum $W_H(0)/W_M(0) \approx 1$. The VSQBO amplitude is defined by eq. (24) and its initial phase is $\varphi_0 = \pi$. a) the latitude $\theta = \theta_2^{(N)}$; b) $\theta = \theta_2^{(S)}$. 